

Solution to Chap.6

**Solution 6.5** Since  $f(x, y)$  is rotationally symmetric,  $g(\ell, \theta) = g(\ell, 0)$ . Hence,

$$\begin{aligned} g(\ell, \theta) &= \int_{L(\ell, \theta)} f(x, y) ds \\ &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-x^2-y^2} \delta(x \cos(\theta) - \ell) dx dy \\ &= \int_{-\infty}^{\infty} e^{-\ell^2-y^2} dy \\ &= e^{-\ell^2} \int_{-\infty}^{\infty} e^{-y^2} dy. \end{aligned}$$

Since

$$\frac{1}{\sqrt{2\pi\sigma^2}} \int_{-\infty}^{\infty} e^{-x^2/\sigma^2} dx = 1,$$

then,

$$\int_{-\infty}^{\infty} e^{-y^2} dy = \sqrt{\pi}.$$

Thus,

$$g(\ell, \theta) = \sqrt{\pi} e^{-\ell^2}.$$

**Solution 6.9** We start with the convolution integral

$$g(x, y) = \iint_{\xi, \eta} f(\xi, \eta) h(x - \xi, y - \eta) d\xi d\eta.$$

Now, we perform the following steps:

$$\begin{aligned} \mathcal{R}\{g\} &= \int_y \int_x \int_{\eta} \int_{\xi} f(\xi, \eta) h(x - \xi, y - \eta) d\xi d\eta \delta(x \cos \theta + y \sin \theta - \ell) dx dy \\ &= \int_{\eta} \int_{\xi} f(\xi, \eta) \int_y \int_x h(x - \xi, y - \eta) \delta(x \cos \theta + y \sin \theta - \ell) dx dy d\xi d\eta \\ &= \int_{\eta} \int_{\xi} f(\xi, \eta) \int_{y'} \int_{x'} h(x', y') \delta(x' \cos \theta + y' \sin \theta - [\ell - \xi \cos \theta - \eta \sin \theta]) dx' dy' d\xi d\eta \\ &= \int_{\eta} \int_{\xi} f(\xi, \eta) \mathcal{R}\{h\}(\ell - \xi \cos \theta - \eta \sin \theta, \theta) d\xi d\eta \\ &= \int_{\eta} \int_{\xi} f(\xi, \eta) \int_{\ell'} \mathcal{R}\{h\}(\ell - \ell', \theta) \delta(\xi \cos \theta + \eta \sin \theta - \ell') d\ell' d\xi d\eta \\ &= \int_{\ell'} \mathcal{R}\{h\}(\ell - \ell', \theta) \int_{\eta} \int_{\xi} f(\xi, \eta) \delta(\xi \cos \theta + \eta \sin \theta - \ell') d\xi d\eta d\ell' \\ &= \int_{\ell'} \mathcal{R}\{h\}(\ell - \ell', \theta) \mathcal{R}\{f\}(\ell', \theta) d\ell' \\ &= \mathcal{R}\{h\} * \mathcal{R}\{f\} \end{aligned}$$

which was to be proved.

**Solution 6.10**

(a)

$$\begin{aligned}
 g_s(\ell, \theta + \pi/2) &= \iint s(x, y) \delta(x \cos(\theta + \pi/2) + y \sin(\theta + \pi/2) - \ell) dx dy \\
 &= \iint s(x, y) \delta(-x \sin \theta + y \cos \theta - \ell) dx dy \\
 &= \iint s(-v, u) \delta(u \cos \theta + v \sin \theta - \ell) du dv \quad (u = y, v = -x) \\
 &= \iint s(u, v) \delta(u \cos \theta + v \sin \theta - \ell) du dv \\
 &= g_s(\ell, \theta)
 \end{aligned}$$

(b)

$$\begin{aligned}
 g_s(\ell, -\theta) &= \iint s(x, y) \delta(x \cos(-\theta) + y \sin(-\theta) - \ell) dx dy \\
 &= \iint s(x, y) \delta(x \cos \theta - y \sin \theta - \ell) dx dy \\
 &= \iint s(u, -v) \delta(u \cos \theta + v \sin \theta - \ell) du dv \quad (u = x, v = -y) \\
 &= \iint s(u, v) \delta(u \cos \theta + v \sin \theta - \ell) du dv \\
 &= g_s(\ell, \theta)
 \end{aligned}$$

(c) Let

$$\tilde{g}_s(\ell, \theta) = \begin{cases} g_s(\ell, \theta) & 0 \leq \theta < \frac{\pi}{4} \\ g_s\left(\ell, \frac{\pi}{2} - \theta\right) & \frac{\pi}{4} \leq \theta < \frac{\pi}{2} \end{cases},$$

which covers  $0 \leq \theta < \pi/2$ . Then

$$g_s(\ell, \theta) = \begin{cases} \tilde{g}_s(\ell, \theta) & 0 \leq \theta < \frac{\pi}{2} \\ \tilde{g}_s\left(\ell, \theta - \frac{\pi}{2}\right) & \frac{\pi}{2} \leq \theta \leq \pi \end{cases},$$

covers  $0 \leq \theta < \pi$ .

(d) See Figure S6.1.

(e) See Figure S6.2.

From simple rotations, we have  $\ell_1 = -\cos \theta - \sin \theta$ ,  $\ell_2 = -\cos \theta + \sin \theta$ ,  $\ell_3 = \cos \theta - \sin \theta$ , and  $\ell_4 = \cos \theta + \sin \theta$ .

By similar triangles, we have:

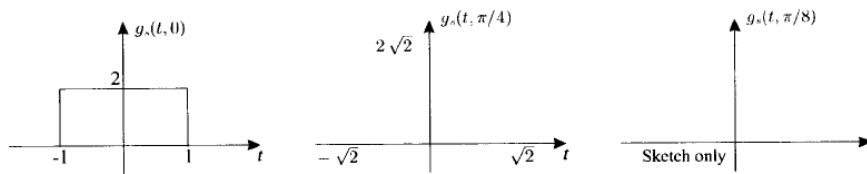


Figure S6.1 Sketch of projections at different angles. [Problem 6.10(d)]

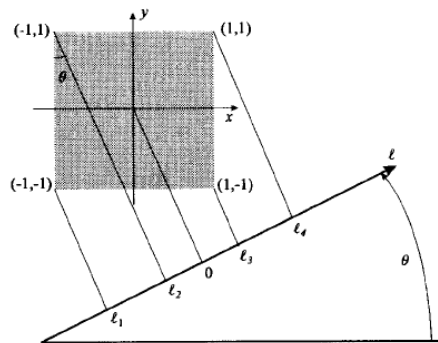


Figure S6.2 [Problem 6.10(e)]

$$\begin{aligned}
 \underline{l_1 \leq l \leq l_2}: \quad \frac{l - l_1}{l_2 - l_1} &= \frac{g(l, \theta)}{\left(\frac{2}{\cos \theta}\right)} \\
 \Rightarrow g(l, \theta) &= \frac{2}{\cos \theta} \left( \frac{l + \cos \theta + \sin \theta}{2 \sin \theta} \right) \\
 &= \frac{l + \cos \theta + \sin \theta}{\cos \theta \sin \theta} \\
 \underline{l_2 \leq l \leq l_3}: \quad g(l, \theta) &= \frac{2}{\cos \theta} \\
 \underline{l_3 \leq l \leq l_4}: \quad \frac{l_4 - l}{l_4 - l_3} &= \frac{g(l, \theta)}{\left(\frac{2}{\cos \theta}\right)} \\
 \Rightarrow g(l, \theta) &= \frac{\cos \theta + \sin \theta - l}{\cos \theta \sin \theta} .
 \end{aligned}$$

and  $g(l, \theta) = 0$  elsewhere.

### Solution 6.13

(a)

$$g(l, 60) = \begin{cases} \sqrt{3}\mu(a/2 + l) & -a/2 \leq l \leq 0 \\ \sqrt{3}\mu(a/2 - l) & 0 \leq l \leq a/2 \\ 0 & \text{otherwise} \end{cases}$$

The projection  $g(\ell, 60^\circ)$  is shown in Figure S6.9.

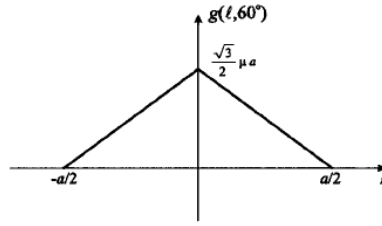


Figure S6.9 [Problem 6.13(a)]

(b)

$$\begin{aligned}
 b_{60^\circ} \left( 0, \frac{a}{4} \right) &= g \left( 0 \cos 60^\circ + \frac{a}{4} \sin 60^\circ, 60^\circ \right) \\
 &= g \left( \frac{\sqrt{3}a}{8}, 60^\circ \right) \\
 &= \sqrt{3}\mu \left( \frac{a}{2} - \frac{\sqrt{3}a}{8} \right) \\
 &= \frac{\sqrt{3}\mu a}{8} (4 - \sqrt{3}).
 \end{aligned}$$

(c) Let function  $f(t)$  be defined by scaling a rect function as:

$$f(t) = \begin{cases} K, & -\frac{p}{2} \leq t \leq \frac{p}{2} \\ 0, & \text{otherwise} \end{cases}.$$

The convolution  $f(t) * f(t)$  is given by:

$$f(t) * f(t) = \begin{cases} K^2(p+t), & -p \leq t \leq 0 \\ K^2(p-t), & 0 \leq t \leq p \\ 0, & \text{otherwise} \end{cases}.$$

Comparing with  $g(\ell, 60^\circ)$ , we see that it is a convolution of function  $f(\ell)$  with itself, where  $K^2 = \sqrt{3}\mu$ , and  $p = a/2$ :

$$g(\ell, 60^\circ) = \sqrt{\sqrt{3}\mu} \operatorname{rect} \left( \frac{2\ell}{a} \right) * \sqrt{\sqrt{3}\mu} \operatorname{rect} \left( \frac{2\ell}{a} \right).$$

By the projection slice theorem, we get  $F(\varrho \cos \theta, \varrho \sin \theta) = G(\varrho, \theta) = \mathcal{F}\{g(\ell, \theta)\}$ . Since  $g(\ell, 60^\circ)$  is expressed as a convolution of a rect function with itself, we have:

$$\begin{aligned}
 F\{g(\ell, 60^\circ)\} &= [F\{\sqrt{\sqrt{3}\mu} \cdot \operatorname{rect}(\frac{2\ell}{a})\}]^2 \\
 &= \frac{\sqrt{3}\mu a^2}{4} \operatorname{sinc}^2 \left( \frac{a\varrho}{2} \right).
 \end{aligned}$$

Therefore,

$$F(\varrho \cos 60^\circ, \varrho \sin 60^\circ) = \frac{\sqrt{3}\mu a^2}{4} \operatorname{sinc}^2 \left( \frac{a\varrho}{2} \right).$$

**Solution 6.17**

(a)  $b_\theta(x, y) = g(x \cos \theta + y \sin \theta, \theta)$ .

(b)

$$\begin{aligned} \ell &= x \cos \theta + y \sin \theta = 1 \cos 30^\circ + 2 \sin 30^\circ \\ &= 0.866 + 1 = 1.866. \end{aligned}$$

$$b_{30^\circ}(1, 2) = g(1.866, 30^\circ) \approx 0.155.$$

(c) No, because  $g(\ell, 30^\circ)$  does not say anything about  $g(\ell, 45^\circ)$ .

(d)  $210^\circ = 30^\circ + 180^\circ$ . Thus, this is the “opposite” projection, and

$$g(\ell, 210^\circ) = g(-\ell, 30^\circ) = 0.155.$$

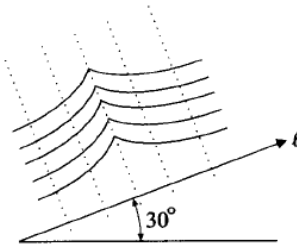
(e) See Figure S6.10. The image always has the same value along the lines with a slope of  $\tan 120^\circ = -\sqrt{3}/3$ .

(f) No, because to determine  $b_{30^\circ}(1, 2)$ , we need  $\ell = 1.866$  as shown in (b), which is not an integer.

An approximate value might be to choose  $\ell = 2$ , which yields 0.135.

(g)  $\ell = 2 \times 0.866 + 1 \times 0.5 = 2.232$ , which is again not an integer. Thus, the exact value still cannot be determined.

For approximation, still  $\ell = 2$ , and the approximate value is still 0.135.



**Figure S6.10** [Problem 6.17(c)]

**Solution 6.19** Beam width  $W$  has the effect of convolving projection with  $\text{rect}(\frac{\ell}{W})$ .

Ignoring sampling (at first) and pretending that we don't know about the distortion, CBP yields

$$\tilde{f}(x, y) = \int_0^\pi \left[ (g_\theta(\ell) * \text{rect}\left(\frac{\ell}{W}\right) * c(\ell)) \right]_{\ell=x \cos \theta + y \sin \theta} d\theta$$

or

$$\tilde{f}(x, y) = \int_0^\pi \int_{-\infty}^\infty \left[ (g_\theta(\ell) * \text{rect}\left(\frac{\ell}{W}\right) * c(\ell)) \right] \delta(x \cos \theta + y \sin \theta - \ell) d\ell d\theta.$$

We let  $g_\theta(\ell) = \delta(\ell) = 2\text{-D Radon transform of } \delta(x, y)$  to find the impulse response. But  $\delta(\ell) * \text{rect}(\frac{\ell}{W}) * c(\ell) = \text{rect}(\frac{\ell}{W}) * c(\ell)$ . Therefore, the impulse response in the inverse 2-D Radon transform of  $g_\theta(\ell) = \text{rect}(\frac{\ell}{W})$ , or the function  $h(x, y)$  whose 2-D Radon transform is  $\text{rect}(\frac{\ell}{W})$ . The function has support on the disk with diameter  $W$  centered at the origin, but is not constant within, as shown in Figure S6.11.

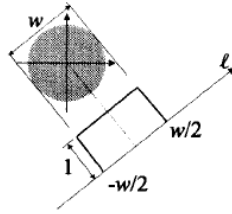


Figure S6.11 [Problem 6.19]

The easiest way to determine  $h(x, y)$  is via the projection-slice theorem.

$$\mathcal{F} \left\{ \text{rect} \left( \frac{\ell}{W} \right) \right\} = |W| \text{sinc}(W\varrho)$$

and since all projections are the same, we conclude that

$$H(\varrho) = |W| \text{sinc}(W\varrho) = |W| \frac{\sin(\pi W\varrho)}{\pi W\varrho}$$

$H(\varrho)$  is the radial part of  $\mathcal{F}(h(x, y))$ . The inverse transform of  $H(\varrho)$  also has circular symmetry and is given by the inverse Hankel transform (Macovski, page 16, for example):

$$h(r) = 2\pi \int_0^\infty H(\varrho) J_0(2\pi\varrho r) \varrho d\varrho$$

where  $J_0$  is the Bessel function of order 0.

From tables (Bracewell for example) we find

$$\frac{\text{rect} \left( \frac{r}{2a} \right)}{\sqrt{a^2 - r^2}} \longleftrightarrow \frac{\sin(2\pi a\varrho)}{\varrho}$$

Hence

$$h(r) = \left( \frac{1}{\pi W} \right) \frac{\text{rect} \left( \frac{r}{W} \right)}{\sqrt{\left( \frac{W}{2} \right)^2 - r^2}}$$

which gives

$$h(x, y) = \left( \frac{1}{\pi W} \right) \frac{\text{rect} \left( \frac{\sqrt{x^2 + y^2}}{W} \right)}{\sqrt{\left( \frac{W}{2} \right)^2 - (x^2 + y^2)}}$$

Finally, we conclude that

$$\tilde{f}(x, y) = f(x, y) * h(x, y).$$

Additional comments:  $h(x, y)$  is a low-pass filter since its Fourier transform decays as a sinc in  $\varrho$ . Hence,  $\tilde{f}(x, y)$  is a blurred version of  $f(x, y)$  as expected. However,  $h(x, y)$  has finite support, so that the blurring is strictly local—in fact contributions occur only from over the disk of radius  $\frac{W}{2}$ .

But  $h(x, y)$  approaches  $\infty$  asymptotically at  $r = \frac{W}{2}$ , which means that the contribution to blurring at exactly

the radius  $\frac{W}{2}$  can be very strong and one might expect to see circular artifacts of radius  $\frac{W}{2}$  near bright point objects.

In a real system CBP is done for sampled data. The convolution  $g_\theta(l) * \text{rec}\left(\frac{l}{w}\right)$  is continuous convolution, however, followed by discrete sampling. Therefore we might write

$$\tilde{f}(x, y) = \frac{\pi}{M} \sum_{j=1}^M T \sum_{i=1}^N [g_\theta(s) * \text{rec}\left(\frac{s}{w}\right)]_{x=iT} C(x \cos \theta_j + y \sin \theta_j - iT)$$

But there is not more we can say analytically about  $\hat{f}(x, y)$  vs.  $f(x, y)$ .

*Note:  $h(x, y) = H^{-1}\{S(\rho)\}$ .*

**Solution 6.20** Assume a rectangularly windowed rho-filter is used. The SNR can be computed using Eq. (6.74). By assumption,  $M = 100$ ,  $C = 0.05$ ,  $\bar{\mu} = 0.15 \text{ cm}^{-1}$ . Since the detectors are touching each other,  $k = 1$ . Since the cylinder has a diameter of 20cm, and the detector dimension is 2.0mm  $\times$  2.0mm, the number of measurements per projection is

$$\frac{20 \text{ cm}}{2 \text{ mm}} = 100.$$

Hence,

$$0.1R/\text{projection} \times \frac{1}{100} \text{projection/measurement} = 0.001R/\text{measurement}.$$

The worst-case intersection length of a beam with the water is 20cm. Therefore, the worst-case  $\bar{N}$  is

$$\begin{aligned} \bar{N} &= 2.5 \times 10^{10} \frac{\text{photons}}{\text{cm}^2 \text{R}} \times 0.04 \text{ cm}^2 \times 0.001 \frac{\text{R}}{\text{measurement}} e^{-0.15 \times 20} \\ &\approx 50 \times 10^3 \text{ photons/measurement.} \end{aligned}$$

Thus,

$$\begin{aligned} \text{SNR} &\approx 0.4kC\bar{\mu}\sqrt{\bar{N}Mw} \\ &= 0.4 \times 1 \times 0.05 \times 0.15 \text{ cm}^{-1} \sqrt{50 \times 10^3 \times 100 \times 0.2} \\ &\approx 1.3 \end{aligned}$$

Also,  $\text{SNR} = 20 \log_{10} 1.3 \approx 2.5 \text{ dB}$ , since the SNR is not a power ratio as defined.

Consider a 4x4 image that contains a diagonal line

$$I=[0,0,0,1;0,0,1,0;0,1,0,0;1,0,0,0];$$

- a) determine its projections in the directions: 0, 45, 90, 135 degrees.
- b) determine the backprojected image from each projection;
- c) determine the reconstructed images by using projections in the 0 and 90 degrees only.
- d) determine the reconstructed images by using all projections. Comment on the difference from c).

Solution:

Original image:

$$I = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix}$$

(a) projections in directions:

$$\begin{aligned} 0 \text{ degree:} & \quad 1, 1, 1, 1 \\ 45 \text{ degree:} & \quad 0, 0, 0, 4, 0, 0, 0 \\ 90 \text{ degree:} & \quad 1, 1, 1, 1 \\ 135 \text{ degree:} & \quad 1, 0, 1, 0, 1, 0, 1 \end{aligned}$$

(b) backprojections in each direction:

Each projection should be normalized by the number of pixels along the projection path.

$$b_{0^\circ} = \frac{1}{4} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \end{pmatrix} \quad b_{45^\circ} = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix}$$

$$b_{90^\circ} = \frac{1}{4} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \end{pmatrix} \quad b_{135^\circ} = \begin{pmatrix} 0 & 1/3 & 0 & 1 \\ 1/3 & 0 & 1/3 & 0 \\ 0 & 1/3 & 0 & 1/3 \\ 1 & 0 & 1/3 & 0 \end{pmatrix}$$

$$(c) \quad b_{0^\circ} + b_{90^\circ} = \frac{1}{2} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \end{pmatrix}$$

$$(d) \quad b_{0^\circ} + b_{45^\circ} + b_{90^\circ} + b_{135^\circ} = \begin{pmatrix} 1/2 & 5/6 & 1/2 & 5/2 \\ 5/6 & 1/2 & 11/6 & 1/2 \\ 1/2 & 11/6 & 1/2 & 5/6 \\ 5/2 & 1/2 & 5/6 & 1/2 \end{pmatrix}$$

Using all 4 backprojections, reconstructed image is better to reflect the original image than (c) using 2 backprojections.